Time and Spatial Series and Transforms

- Z- and Fourier transforms
- Gibbs' phenomenon
- Transforms and linear algebra
- Wavelet transforms
 - Reading:
 - Sheriff and Geldart, Chapter 15

Z-Transform

- Consider a discretized record of N readings: $U=\{u_0, u_1, u_2, ..., u_{N-1}\}$. How can we represent this series differently?
- The Z transform simply associates with this time series a *polynomial function*:

$$\{u_i\} \rightarrow U(z) = u_0 + u_1 z + u_2 z^2 + u_3 z^3 + \dots$$

- For example, a 3-sample record of {1,2,5} is represented by a quadratic polynomial: 1+2z+5z².
- In Z-domain, the all-important operation of convolution of time series becomes simple multiplication of their Z-transforms:

$$u_1(t) * u_2(t) \rightarrow U_1(z)U_2(z)$$

Fourier Transform

- To describe a polynomial function of order N-1, it is sufficient to specify its values at N points in the plane of complex variable "z"
- The Discrete Fourier transform is obtained by taking the Z-transform at N points uniformly distributed around a unit circle on the complex plane of z:

$$U(k) = \sum_{m=1}^{N-1} e^{i\frac{2\pi k}{N}m} u(t_m) \qquad k = 0, 1, 2, ..., N-1$$

 Each term (k>0) in the sum above is a periodic function (a combination of sin and cos), with a period of N/k sampling intervals:

$$e^{i\alpha} = \cos(\alpha) + i\sin(\alpha)$$

- Thus, the Fourier transform expresses the signal as a sum of its frequency components,
 - Fourier transform also has the property of the Ztransform regarding convolution

Matrix form of Fourier Transform

 Note that the Fourier transform can be written as matrix multiplication:

$$\begin{pmatrix} U(\omega_1) \\ U(\omega_2) \\ U(\omega_3) \\ \dots \end{pmatrix} = \mathbf{F} \begin{pmatrix} u(t_1) \\ u(t_2) \\ u(t_3) \\ \dots \end{pmatrix}$$

$$\mathbf{F} = egin{bmatrix} e^{i\omega_1t_1} & e^{i\omega_1t_2} & e^{i\omega_1t_3} & \dots \ e^{i\omega_2t_1} & e^{i\omega_2t_2} & e^{i\omega_2t_3} & \dots \ e^{i\omega_3t_1} & e^{i\omega_3t_2} & e^{i\omega_3t_3} & \dots \ \dots & \dots & \dots \end{bmatrix}$$

Inverse:

$$\mathbf{F}^{-1} = \frac{\mathbf{\bar{F}}^{T}}{N} = \frac{1}{N} \begin{bmatrix} e^{-i\omega_{1}t_{1}} & e^{-i\omega_{2}t_{1}} & e^{-i\omega_{3}t_{1}} & \dots \\ e^{-i\omega_{1}t_{2}} & e^{-i\omega_{2}t_{2}} & e^{-i\omega_{3}t_{2}} & \dots \\ e^{-i\omega_{1}t_{3}} & e^{-i\omega_{2}t_{3}} & e^{-i\omega_{3}t_{3}} & \dots \\ \dots & \dots & \ddots \end{bmatrix}$$

Resolution of Fourier Transform

Resolution matrix:

$$\mathbf{R}_F = \mathbf{F}^{-1}\mathbf{F}$$

 If all N frequencies are used to reproduce the Fourier-transformed signal, the recovery is accurate:

$$\mathbf{R}_F = \mathbf{I}$$

 If fewer than N frequencies are used for recovering the signal (Gibbs phenomenon), the resolution is incomplete:

$$\mathbf{R}_F \neq \mathbf{I}$$

Integral Fourier Transform

For continuous time and frequency (infinitesimal sampling interval and infinite recording time), Fourier transform reads:

• Forward:
$$U(\omega) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} dt u(t) e^{i\omega t}$$

• Inverse:
$$u(t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} d\omega U(\omega) e^{-i\omega t}$$

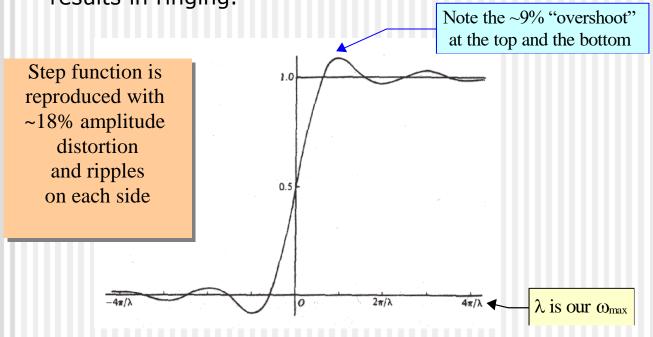
In practice, the bandwidth (and time) is always limited, and so the actual combination of the forward and inverse transforms is rather:

$$u_{\text{band-limited}}(t) = \frac{1}{2\pi} \int_{-\omega_{\text{max}}}^{\omega_{\text{max}}} d\omega \left[\int_{-\infty}^{\infty} d\tau u(\tau) e^{i\omega\tau} \right] e^{-i\omega t}$$

$$u_{\text{band-limited}}(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\tau u(\tau) \left[\int_{-\omega_{\text{max}}}^{\omega_{\text{max}}} d\omega e^{i\omega(t-\tau)} \right]$$

Gibbs' phenomenon

 At a discontinuity, application of the Fourier forward and inverse transform (with a limited bandwidth), results in ringing.

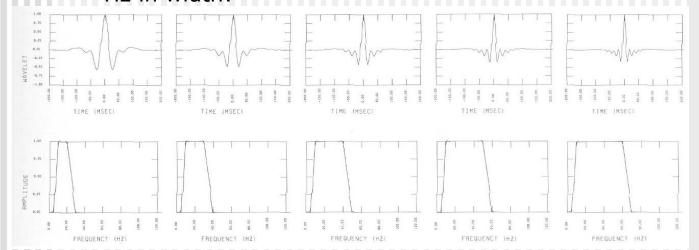


- This is important for constructing time and frequency windows
 - Boxcar windows create ringing at their edges.
 - "Hanning" (cosine) windows are often used to reduce ringing:

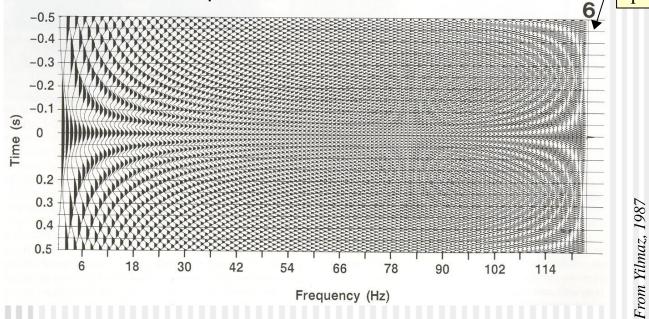
$$H_{\Delta t}(t) = \frac{1}{2} \left(1 - \cos \frac{\pi t}{\Delta t} \right)$$

Spectra of Pulses

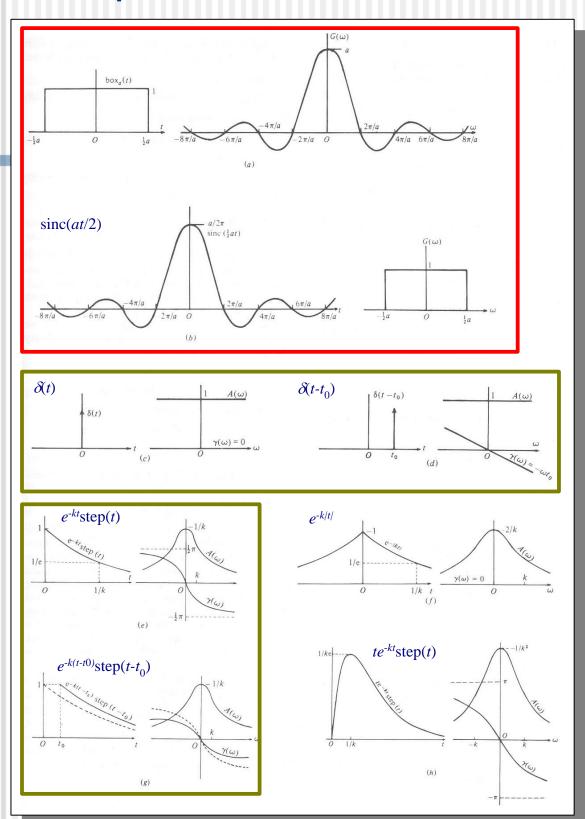
For a pulse of width T s, its spectrum is about 1/T Hz in width:



Equal-amplitude (co)sinusoids from 0 to f_N add up to form a spike:



Sample Fourier Transforms



Wavelet transforms

 Like the inverse Fourier transform,
wavelet decomposition represents the time-domain signal by a combination of wavelets of some desired shapes:

$$\begin{pmatrix} u(t_1) \\ u(t_2) \\ u(t_3) \\ \dots \end{pmatrix} = \begin{bmatrix} f_1(t_1) & f_2(t_1) & f_3(t_1) & \dots \\ f_1(t_2) & f_2(t_2) & f_3(t_2) & \dots \\ f_1(t_3) & f_2(t_3) & f_3(t_3) & \dots \\ \dots & \dots & \dots \end{bmatrix} \begin{pmatrix} a_1 \\ a_2 \\ a_3 \\ \dots \end{pmatrix}$$

Wavelet shapes

Wavelet amplitudes

 Ideally, wavelets should form a complete orthonormal basis:

$$\sum_{k=0}^{N-1} f_i(t_k) f_j(t_k) = \delta_{ij}$$
 although this is not always necessary
$$\underbrace{\sum_{k=0}^{N-1} f_i(t_k) f_j(t_k)}_{\text{exp(...) functions used in Fourier transforms satisfy this property}}_{\text{exp(...) functions used in Fourier transforms satisfy this property}}$$

 Usually, functions f(t) represent time-scaled and shifted versions of some "wavelet" W(t)